

# Quantum Mechanics GRE Notes

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October 20, 2005

Here is some stuff that you should know about quantum mechanics for the GRE.

- *Schrödinger Equation.* General and time-independent:

$$i\hbar \frac{\partial \Psi}{\partial t} = \left( -\frac{\hbar^2}{2m} \nabla^2 + V \right) \Psi$$

$$\begin{aligned} H\psi &= E\psi \\ \Psi &= \psi e^{-iEt/\hbar} \end{aligned} \quad (1)$$

- *Formalism and Operator Algebra.* This is probably actually the easiest topic for people who have taken 143a. Some equations of use are:

$$\langle \alpha | \beta \rangle = \langle \beta | \alpha \rangle^*$$

$$\hat{O} = \hat{O}^\dagger \text{ if } \hat{O} \text{ is Hermitian}$$

$$\begin{aligned} \langle \alpha | \hat{O} | \beta \rangle &= \beta \langle \alpha | \beta \rangle \text{ if } |\beta\rangle \text{ is an eigenstate of } \hat{O} \\ &= \int_{-\infty}^{\infty} dx \alpha^*(x) \hat{O}(x) \beta(x) \text{ in 1D} \\ &= A^\dagger \times O \times B \text{ as matrices} \end{aligned}$$

$$\begin{aligned} [\hat{A}, \hat{B}] &= \hat{A}\hat{B} - \hat{B}\hat{A} \\ \therefore [\hat{A}, \hat{B}] &= -[\hat{B}, \hat{A}] \\ [\hat{A}\hat{B}, \hat{C}] &= [\hat{A}, \hat{C}]\hat{B} + \hat{A}[\hat{B}, \hat{C}] \end{aligned} \quad (2)$$

A Hermitian matrix has real eigenvalues; observables thus must be represented by Hermitian operators.

- *Position and momentum.*

$$\begin{aligned} \hat{p} &= -i\hbar \nabla \\ [\hat{x}, \hat{p}] &= i\hbar \\ [\hat{f}(x), \hat{p}] &= i\hbar \frac{df}{dx} \end{aligned} \quad (3)$$

- *De Broglie equation.* This one always trips me up even though it's easy:

$$\begin{aligned}\lambda &= h/p \\ &= h/\sqrt{2mk_B T}\end{aligned}\tag{4}$$

if the particle's energy is thermal. Fun fact: the De Broglie wavelength of thermal neutrons is about 1.78 angstroms.

- *Emitted photons.* Yeah. The name of the game:

$$\begin{aligned}E &= \frac{-\mu e^4 Z^2}{2h^2 n^2} \\ &= (-13.61\text{eV}) \frac{Z^2}{n^2} \text{ for } \mu = m_e \\ \therefore \lambda^{-1} &\propto Z^2 \left( \frac{1}{n^2} - \frac{1}{n'^2} \right)\end{aligned}\tag{5}$$

This means the key to their beloved positronium problems is that  $\mu \rightarrow \mu/2$  since the reduced mass is now  $(m_e^{-1} + m_e^{-1})^{-1}$ . So basically you can do the problem as if it were hydrogen with  $13.6 \text{ eV} \rightarrow 6.8 \text{ eV}$ .

- *Hydrogen Ground State.* A good thing to know:

$$\psi(r) = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}\tag{6}$$

- *Infinite Square Well.* Straightforward:

$$\begin{aligned}\psi_n &= \frac{1}{\sqrt{a}} \sin \frac{n\pi x}{a} \\ E_n &= \frac{\pi^2 n^2 \hbar^2}{2ma^2}\end{aligned}\tag{7}$$

Pay attention to how  $E_n$  scales with  $n$ . Also,  $n \geq 1$  of course.

- *Singlet and triplet states.* These are useful from two perspectives. The first point of view is taking them as eigenstates of particle exchange in a two-particle system. The second is ways of adding two spin- $\frac{1}{2}$  particles to get one spin-1 particle. From the first perspective, the singlet state has eigenvalue -1; from the second, it has  $s = m = 0$ :

$$|0\ 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)\tag{8}$$

The triplet states have eigenvalue 1 with regards to parity and have  $s = 1$  and  $m \in \{-1, 0, 1\}$ :

$$\begin{aligned}|1\ 1\rangle &= |\uparrow\uparrow\rangle \\ |1\ 0\rangle &= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \\ |1\ -1\rangle &= |\downarrow\downarrow\rangle\end{aligned}\tag{9}$$

- *Selection rules.* They seem to like making you remember these.

$$\begin{aligned}\Delta l &= \pm 1 \\ \Delta m_l &= \pm 1 \text{ or } 0 \\ \Delta j &= \pm 1 \text{ or } 0\end{aligned}\tag{10}$$

Here  $l$  is the orbital angular momentum;  $m_l$  is the  $z$  component of  $l$ ;  $j = l + s$  is the total angular momentum. (I'm not exactly sure how you can get  $\Delta j = 0$  but I read it on the internet so it must be true!) Typically  $\Delta n = \pm 1$  but there is no reason that this must hold. Just remember that the emitted photon is spin 1 and so its angular momentum can be described as one of  $|0\ 0\rangle$ ,  $|1\ 1\rangle$ ,  $|1\ 0\rangle$ , or  $|1\ -1\rangle$ .

The phrase "electric dipole radiation" implies that  $\Delta l = 0$  is not physically permissible, even though it is allowed mathematically. Electric dipole transitions are "allowed", while magnetic dipole or electric quadrupole transitions are "forbidden" (but do actually occur).

- *Observed states and expectation values.* The usual format of such a problem is

$$|\psi\rangle = \sqrt{p_1}|k_1\rangle + \sqrt{p_2}|k_2\rangle + \sqrt{p_3}|k_3\rangle + \dots\tag{11}$$

The expectation value is  $p_1 k_1 + p_2 k_2 + p_3 k_3 + \dots$  if  $k_i$  are the eigenvalues associated with eigenkets  $|k_i\rangle$ , and the probability of being in a state  $|k_i\rangle$  is just  $p_i$ . This is all assuming the  $|k_i\rangle$  are eigenkets of whatever operator you're finding the expectation value of.

- *Heisenberg relation.* The general version:

$$\Delta A \Delta B \geq \frac{1}{2} |\langle [A, B] \rangle|\tag{12}$$

In particular,

$$\begin{aligned}\Delta x \Delta p &\geq \hbar/2 \\ \Delta E \Delta t &\geq \hbar/2\end{aligned}\tag{13}$$

- *Time Independent First-Order Perturbation Theory.* Given a perturbation  $H'$  to some original Hamiltonian with eigenstates  $|\psi_n^0\rangle$  and energies  $E_n^0$ , the perturbed energies and states are:

$$\begin{aligned}E_n^1 &= E_n^0 + \langle \psi_n^0 | H' | \psi_n^0 \rangle \\ \psi_n^1 &= \psi_n^0 + \sum_{m \neq n} \frac{\langle \psi_m^0 | H' | \psi_n^0 \rangle}{E_n^0 - E_m^0} \psi_m^0\end{aligned}\tag{14}$$

Recall that

$$\langle \psi_n | \hat{O} | \psi_m \rangle = \int_{-\infty}^{\infty} \psi_n^*(x) \hat{O}(x) \psi_m(x) dx\tag{15}$$

in one dimension. (You will often be perturbing the infinite square well.)

- *Fundamental Understanding of the Concepts of Quantum Mechanics.* Yeah. Whatever.